

Generating Random Gaussian States

References

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A Quick Review on Gaussian States

■ Gaussian States are fully characterized by it's Covariance Matrix (CM); i.e. expectations on position \hat{x} and momentum \hat{p} follow

a Gaussian distribution. Namely,
 given the density matrix ρ
 and the expectation $\langle \cdot \rangle = \text{tr}(\rho \cdot)$,
 we define the CM as:

$$S_\rho := \frac{1}{2} \langle \{ \delta x_i, \delta x_j^T \} \rangle, \text{ with}$$

$$X = (x_1, x_2, \dots, x_{2n}) \\
 = (\hat{q}_1, \hat{p}_1, \hat{q}_2, \hat{p}_2, \dots, \hat{q}_n, \hat{p}_n),$$

with \hat{q}_i, \hat{p}_i being the ~~with~~
 quadratures of the i -th
 bosonic mode.

▣ The canonical commutation
 relations can be expressed very
 concisely as:

$$[x_i, x_j] = i \delta_{ij}^{(J)}, \text{ with}$$

$$J_{2n} := \bigoplus_{i=1}^n \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \text{ being called}$$

the symplectic form. This is central
 in the algebraic structure of

Gaussian states.

→ why do we care about this stuff?

② Physicality of states:

Valid quantum states obey:

$$\{ \Sigma_p + i \Sigma_{2n} \geq 0 \}$$

(This is why we can't just sample from GOE to get random states).

This is basically another way to express the uncertainty principle.

② Simple computation of relevant physical quantities entirely from the CM.

For instance, the purity is:

$$\{ \text{tr}(\rho^2) = \frac{1}{\sqrt{\det \Sigma_p}} \}$$

④ A very convenient diagonalization following the underlying algebra. This is Williamson's Theorem. One can write

$$S_g = M \text{diag}(\nu_1, \nu_1, \nu_2, \nu_2, \dots, \nu_N, \nu_N) M^T$$

Here:

- M is a symplectic transformation, it obeys $M^T J M = J$, i.e. it leaves the symplectic form invariant

- $\{\nu_i\}$ are the symplectic eigenvalues, defined as:

$$\nu_i = \text{eig}(iJ \text{Im} S_g) \quad (\text{s.t. } \nu_i > 0)$$

Note: these eigenvalues come in positive negative pairs $\{\nu_1, -\nu_1, \nu_2, -\nu_2, \dots\}$.

Generating the Random State

Previous technique = by

Fukuda & König, and Serafini.

Major Challenge:

The group

$Sp(2n, \mathbb{R}) = \{M \in 2n \times 2n \text{ real sym. matrices} :$

$$M J_{2n} M^T = J_{2n}\}$$

is not compact, sampling is hard! (Energy is unbounded for Gaussian states).

They perform some random (subset) of symplectic rotations on a pure diagonal matrix after imposing an upper bound on energies.

What do we do?

Strategy: start with arbitrary

GOE matrix G .

This likely won't satisfy

$G + iJ_{2n} \geq 0$. Solution: transfer

it to $G \rightarrow G + \lambda I$, i.e. shift

it by a multiple of the identity in such a way that we satisfy this bonafide

relation.

How? This can be cast as

a semi-definite programming

problem.

$$\min \|G - H\|$$

real $2n \times 2n$

$$(*) \text{ s.t. } H \in \mathcal{M}_{2n}^{\text{sym}}(\mathbb{R}) \uparrow \text{ sym. matrices}$$

$$H \succeq iJ_{2n}$$

The norm here is, in principle, arbitrary. We'll choose the operator norm, with $\|A\|_{\infty} = \max_i |\lambda_i|$

Introduce using the notion

of Rayleigh coefficients, finding the minimum / maximum eigenvalue can also be cast as an SDP problem!

$\lambda(\max)$ is $\frac{x^* M x}{x^* x}$ is maximum for $x = v_{\max}$
 $\lambda(\min)$ is $\frac{x^* M x}{x^* x}$ is minimum for $x = v_{\min}$

$$\text{Maximize } \text{tr}(M Z)$$

$$\text{subject to } \langle I, Z \rangle = 1$$

$$Z \succeq 0$$

Or, its dual:

$$\text{Minimize } y$$

$$\text{subject to } y I - M \succeq 0$$

SDP for operator norm $\|A\|_\infty$, chapter 2 of the SDP book.

Thus, one minimizes the scalar $t = \|G - H\|_\infty$, choosing the operator

$$\begin{aligned} & \|G - H\|_\infty \\ & \min t \\ & \text{s.t. } H \in \mathcal{M}_{2n}^{\text{sa}}(\mathbb{R}), t \in \mathbb{R} \\ & H \succeq iJ_{2n} \end{aligned}$$

$$-tI_{2n} \leq G - H \leq tI_{2n}, \text{ with dual}$$

$$\begin{aligned} & \max \langle Z, iJ_{2n} - G \rangle \\ & \text{s.t. } Z \in \mathcal{M}_{2n}^{\text{sa}}(\mathbb{C}) \\ & Z \succeq 0 \\ & \text{Hermitian} \\ & \underline{\text{Tr } Z \leq 1} \end{aligned}$$

Similar to largest eigenvalue problem, but w/ absolute value instead.

with solution

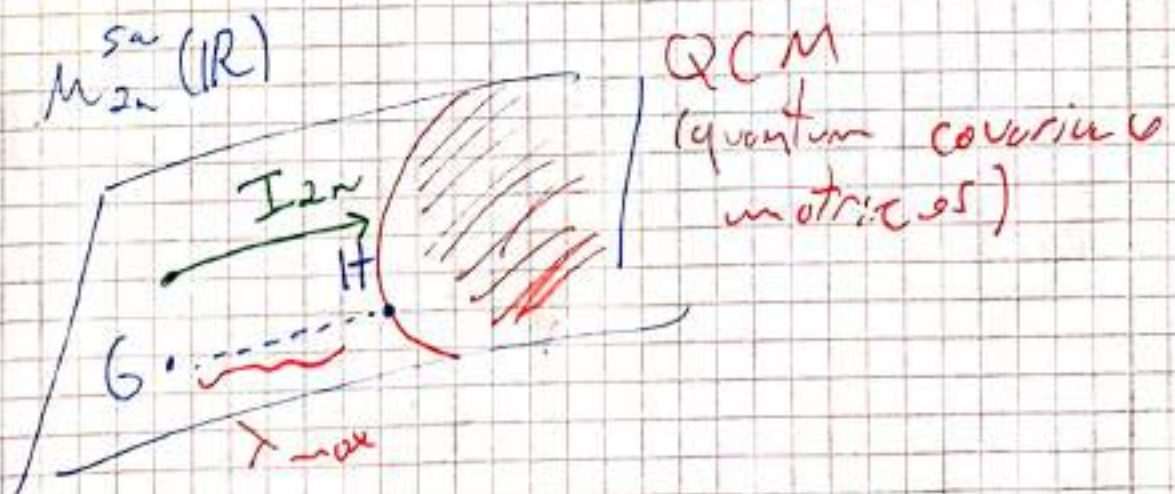
$$\lambda(G) = \max(0, \lambda_{\max}[iJ_{2n} - G])$$

where $\lambda_{\max}[\dots]$ denotes the largest eigenvalue of a matrix.

$$\text{And } H = G + \lambda(G)I_{2n}$$

Note that this last SDP is just the SDP for the matrix $iJ_{2n} - G$, with an inequality (instead of equality) on the trace.

And overall, in the solution we simply shift the (possibly non-physical) matrix G by $\lambda(G)I$, bringing it into the convex hull of valid Gaussian states.



Random Quantum Covariance
Matrices (RQCM)

Let $G \in \text{GOE}(2n, \sigma)$.
we deviate by $\text{RQCM}(2n, \sigma)$
the ensemble:

$$S_G := G + \lambda_{\max} [iJ_{2n} - G] \cdot I_{2n}$$

which are all valid CMs.

Proof:

$$\lambda_{\min} [S_G - iJ_{2n}]$$
$$= \lambda_{\min} [G - iJ_{2n}] + \lambda_{\max} [iJ_{2n} - G]$$

≥ 0 , because $\lambda_{\min} [-x] = -\lambda_{\max} [x]$.

Note that those matrices in
the RQCM ensemble have
 $K(2n) = O(2n) \cap Sp(2n)$ symmetry
ortho-symplectic group

That is,

$$S_{UGUT} = U S_G U^T$$

Proof: $S_{UGUT} = U G U^T + \lambda_{\max} [I_{2n} - U G U^T] I_{2n}$

use orthogonality $U^T U = I$
 $= U G U^T + \lambda_{\max} [U (I_{2n} U^T - G) U^T] I_{2n}$

use simplification in variable $U^T U = I_{2n}$
 $= U G U^T + \lambda_{\max} [I_{2n} - G] I_{2n}$

$$= U G U^T + \lambda_{\max} [I_{2n} - G] I_{2n}$$

orthogonality $U^T U = I_{2n}$
 $= U (G + \lambda_{\max} [I_{2n} - G]) U^T$

$$= U S_G U^T$$

Remarks: for the derivation of the SDP problems, check the Example 2.3 in the book, page 57; as well as examples 2.1 and 2.4.

Interlude: A positive $2n \times 2n$ matrix S_p has decomposition

$$S_p = M \text{diag}(v_1, v_1, v_2, v_2, \dots, v_n, v_n) M^T,$$

or $n \times n$ form.

We get the set $\{v_i\}$ as the positive eigenvalues of $i \Omega M$,

since

$$i \Omega M = B (I_n \otimes -\sigma_2) B^{-1}, \text{ with}$$

$$B = M^T (I_n \otimes U_2), \text{ w/ } U_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ i & -i \end{pmatrix}.$$

That is,

$$D_{U \otimes -\sigma_2} = \text{diag}(-v_1, v_1, -v_2, v_2, \dots);$$

we pick the positive signals of $i \Omega M$, and

$$B = M^T (I_n \otimes U_2) = J_{2n} M J_{2n}^T (I_n \otimes U_2)$$

